

EEE321 Electromagnetic Fields and Waves

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(8th Week)

Outline

- Lorentz Force
- Magnetic Force on a Current-Carrying Wire
- Magnetic Force between two Closed Circuit
- Magnetic Force on a Volume Distribution of Current

The Lorentz Force

- Experiments Show that the force exerted on a particle of charge Q moving in a vacuum at an instantaneous velocity \mathbf{v} in a region where exist both an electric field and magnetic field is

$$\mathbf{F} = Q(\mathbf{E} + \mathbf{v} \times \mathbf{B}). \quad (22-1)$$

- This equation is valid even if \mathbf{v} approaches the speed of light. \mathbf{E} , \mathbf{B} and \mathbf{v} can be time dependent but all concern the same reference frame
- Magnetic force ($Q\mathbf{v} \times \mathbf{B}$) is magnetic force and is perpendicular to \mathbf{v} . Therefore
 - it can change the direction of \mathbf{v}
 - it can not alter its magnitude
 - It can not alter the kinetic energy

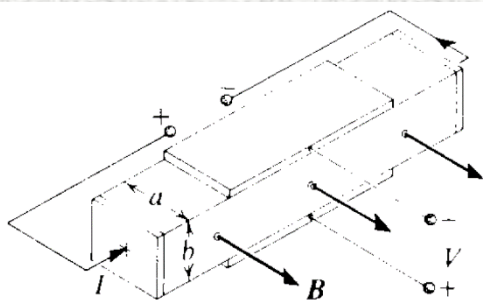
Example: The Hall Effect

- In a bar of conducting material subjected to electric and magnetic field as in figure (a) and (b)
- They also drift sideways because of $Qv \times B$ force
- As a result there appears a potential difference between the upper and lower electrodes
- This is called the **Hall Effect**

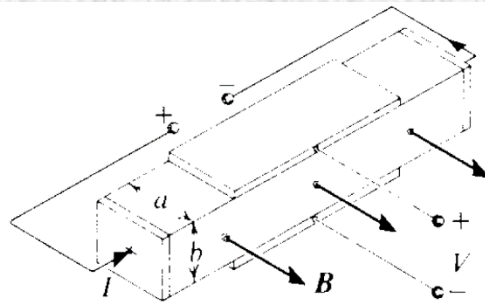
$$QE_t = \frac{QV}{b} = QvB, \quad (22-2)$$

$$V = vBb. \quad (22-3)$$

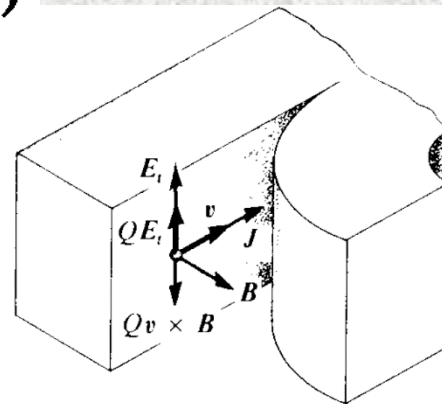
$$V = E_{\text{long}} \mu B b. \quad (22-4)$$



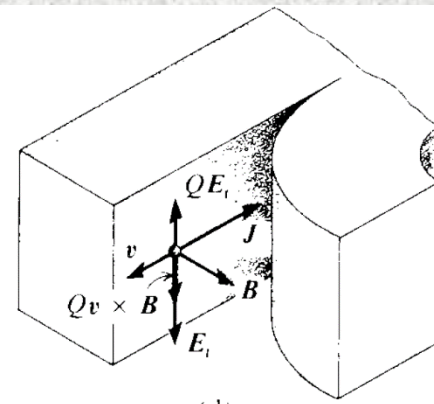
(a)



(b)



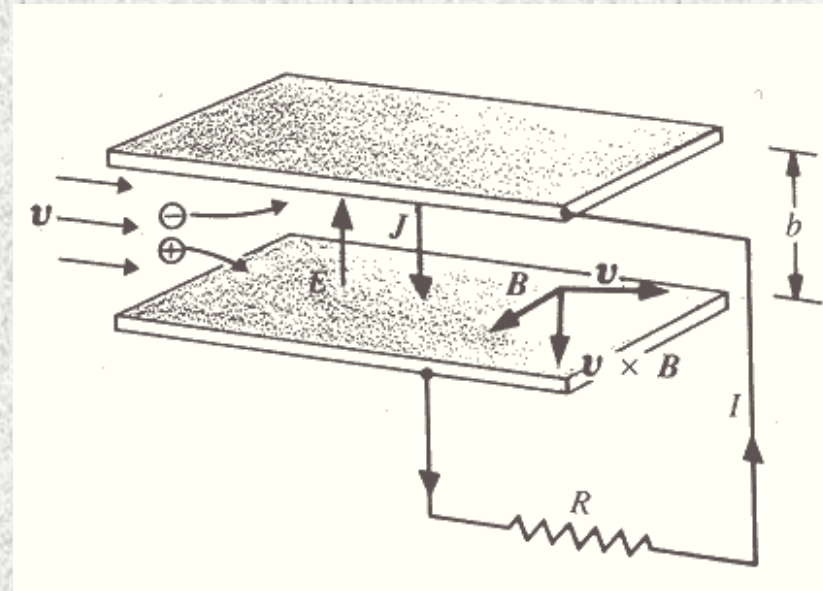
(c)



(d)

Example: Magnetohydrodynamic Generator

- Magnetohydrodynamic (MHD) generator is a large scale application of the hall effect.
- It converts part of the kinetic energy of a hot gas directly into electric energy
- A hot gas enters on the left at a velocity of the order of 1000 m/s
- It contains salt such as (K_2CO_3) that ionizes readily at high temp.
- The temp. approaches 3000 K⁰
- Conductivity is about 100 S/m
- Positive ions curve downward and electrons upward
- The resulting current I flows through the load resistance R
- The advantage of MHD generator is that comprises no moving parts except for the gas



- The function of magnetic force is to compel the positive particules to go to the positive electrode and the negative particules to go to negative electrode.
- Under the action of the forces shown in figure, both move uphill and slow down.
- The Lorentz force acts on a charge Q as if the electric field strength were $\mathbf{E} + \mathbf{v} \times \mathbf{B}$. So for the gas conductivity σ ,

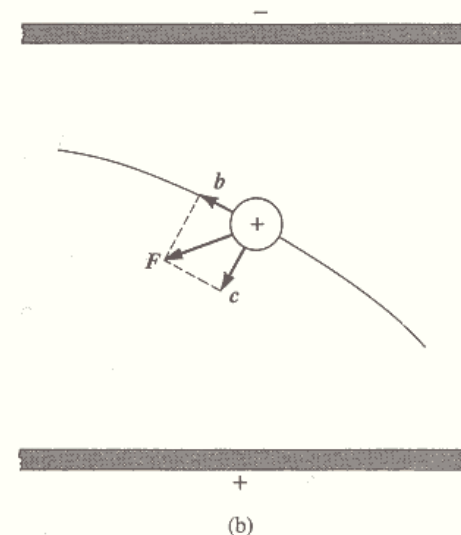
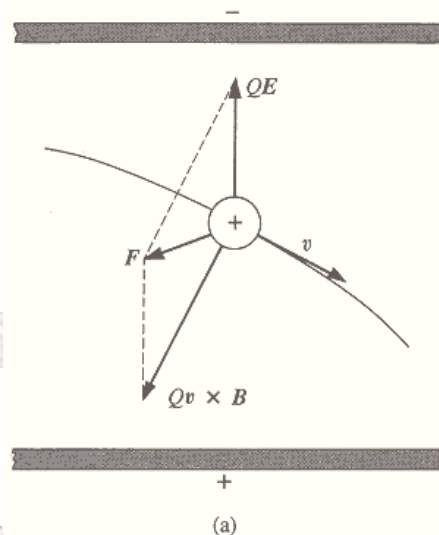
$$|\mathbf{J}| = |\sigma(\mathbf{E} + \mathbf{v} \times \mathbf{B})| = \sigma(vB - E) = \sigma\left(vB - \frac{V}{b}\right), \quad (22-6)$$

$$I = \mathcal{A}J = \mathcal{A}\sigma\left(vB - \frac{IR}{b}\right),$$

$$I = \frac{vBb}{b/(\sigma\mathcal{A}) + R}.$$

- The output voltage is

$$V = IR = vBb - \frac{Ib}{\sigma\mathcal{A}}.$$



- Let us define the efficiency as

$$\mathcal{E} = \frac{\text{Joule losses in } R}{\text{Joule losses in } R + \text{Joule losses in the gas}} \quad (22-10)$$

$$= \frac{I^2 R}{I^2 R + I^2 b / (\sigma \mathcal{A})} = \frac{R}{R + b / (\sigma \mathcal{A})}, \quad (22-11)$$

- Form the above expression for I and for IR ,

$$\mathcal{E} = R \frac{I}{v B b} = \frac{IR}{v B b} = 1 - \frac{I}{\sigma \mathcal{A} v B}. \quad (22-12)$$

- The efficiency is equal to unity when $I=0$ or when $R \rightarrow \infty$.
- The efficiency is equal to zero when $I=\sigma \mathcal{A} v B$ or $R=0$, $V=0$

Magnetic Force on a Current Carrying Wire

- A stationary wire of cross section A carries a current I in a region where there exist a magnetic field \mathbf{B} originating elsewhere
- The wire contains N charge carriers per cubic meter drifting at a velocity \mathbf{v} , each one of charge Q .
- An element of length $d\mathbf{l}$ of a wire contains $ANd\mathbf{l}$ charge carriers. Then the magnetic force on $d\mathbf{l}$ is

$$d\mathbf{F} = AN d\mathbf{l} Q\mathbf{v} \times \mathbf{B} = (ANQ\mathbf{v}) d\mathbf{l} \times \mathbf{B} = I d\mathbf{l} \times \mathbf{B}, \quad (22-13)$$

- The magnetic force per unit length is therefore $I\mathbf{x}\mathbf{B}$
- The total magnetic force on a closed circuit C carrying a current I and lying in a magnetic field is

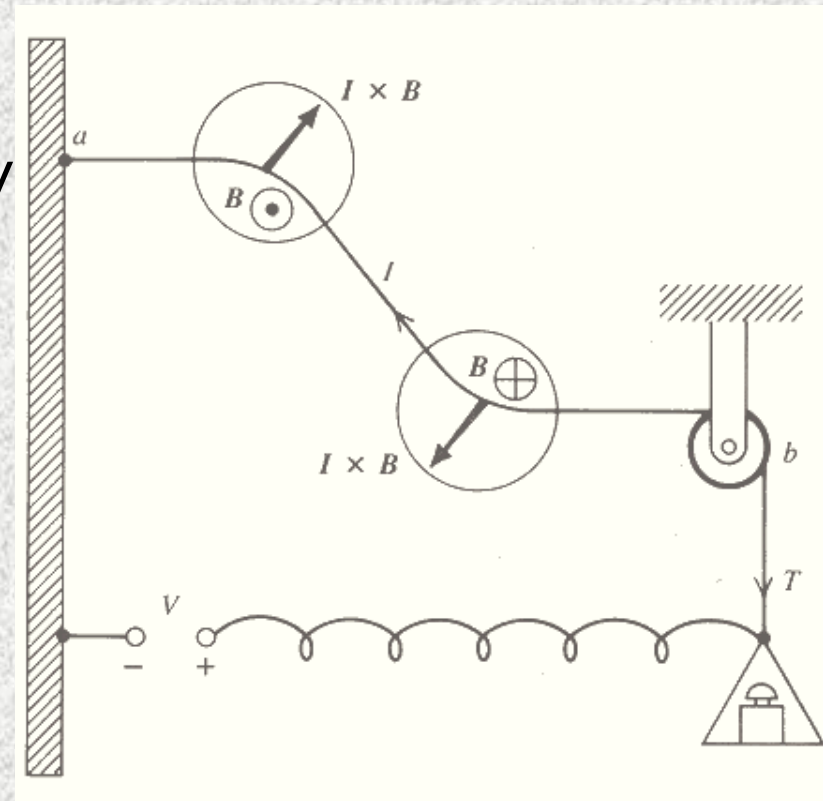
$$\mathbf{F} = I \oint_C d\mathbf{l} \times \mathbf{B}. \quad (22-14)$$

Example: Floating-Wire Hodoscope

- The floating-wire hodoscope is a device that simulates the trajectory of a charged particle in a magnetic field.
- Say a charged particle of mass m , charge Q and velocity \mathbf{v} follows a certain trajectory in going from a point a to a point b in a magnetic field.
- A light wire carrying a current I , fixed at a and going over a pulley at b , in the same magnetic field, as in the figure, will follow that trajectory if

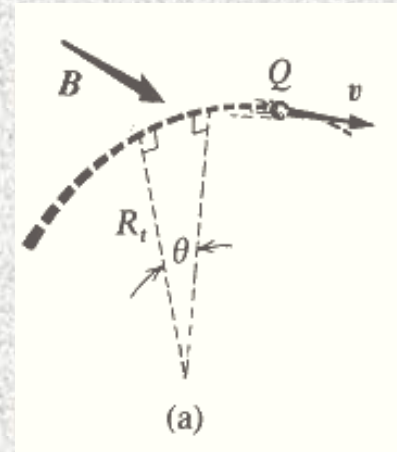
$$\frac{mv}{Q} = \frac{T}{I}, \quad (22-15)$$

- T is the tension in the wire



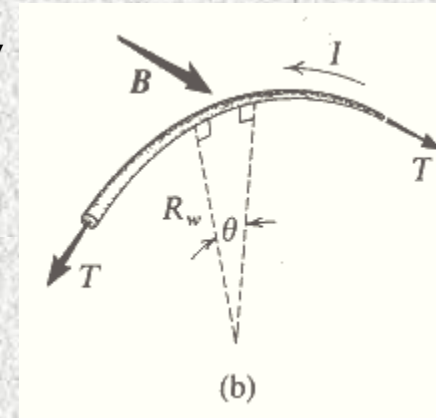
- Suppose the beam is normal to \mathbf{B} as in fig. (a)
- R_t is the radius of curvature of the trajectory

$$QvB = \frac{mv^2}{R_t}, \quad R_t = \frac{mv}{QB}. \quad (22-16)$$



- Suppose the wire is also normal to \mathbf{B} as in fig(b)
- The element dl , with a radius of curvature R_w is in equilibrium if the outward magnetic force $BI dl$ just compensates the inward component of the tension force T :

$$BI dl = 2T \sin \frac{d\theta}{2} = T \frac{dl}{R_w}, \quad R_w = \frac{T}{IB}. \quad (22-17)$$



- The two radii are equal if $\frac{mv}{Q} = \frac{T}{I}. \quad (22-18)$

Magnetic Force between two Closed Circuits

- The magnetic force exerted by a current I_a on a current I_b is given by

$$\begin{aligned} \mathbf{F}_{ab} &= I_b \oint_b d\mathbf{l}_b \times \frac{\mu_0}{4\pi} I_a \oint_a \frac{d\mathbf{l}_a \times \hat{\mathbf{r}}}{r^2}, \\ &= \frac{\mu_0}{4\pi} I_a I_b \oint_a \oint_b d\mathbf{l}_b \times \frac{d\mathbf{l}_a \times \hat{\mathbf{r}}}{r^2}, \end{aligned}$$

- To Show that $\mathbf{F}_{ab} = -\mathbf{F}_{ba}$, first

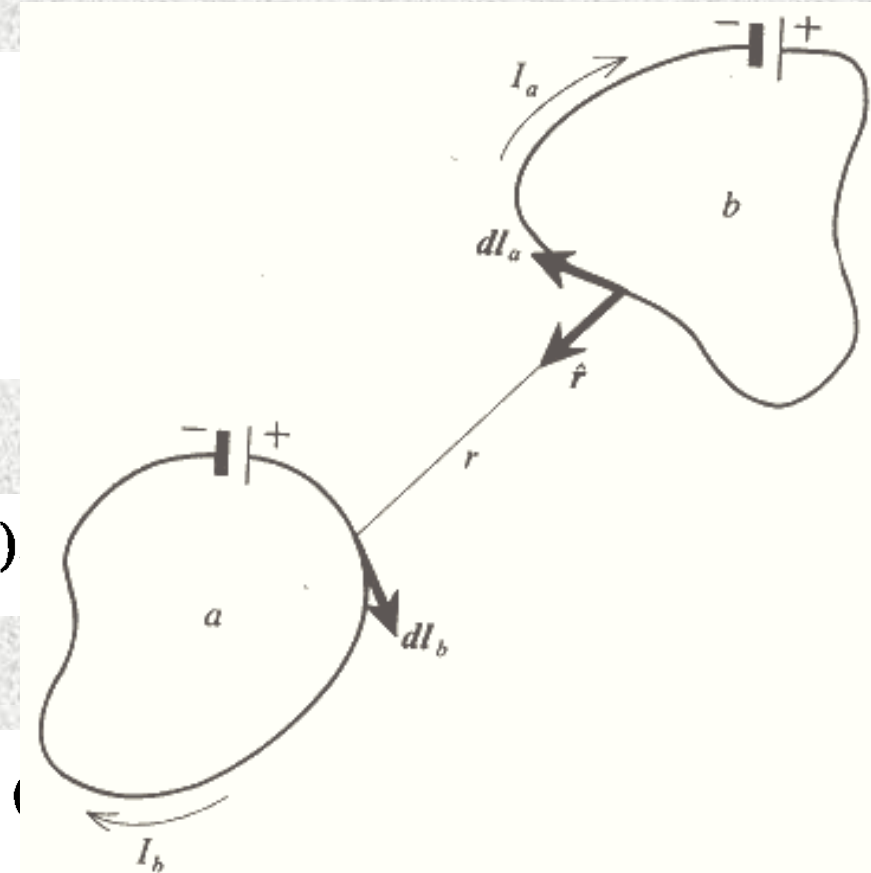
$$d\mathbf{l}_b \times (d\mathbf{l}_a \times \hat{\mathbf{r}}) = d\mathbf{l}_a (d\mathbf{l}_b \cdot \hat{\mathbf{r}}) - \hat{\mathbf{r}} (d\mathbf{l}_a \cdot d\mathbf{l}_b)$$

- Then, rearranging terms,

$$\oint_a \oint_b \frac{d\mathbf{l}_a (d\mathbf{l}_b \cdot \hat{\mathbf{r}})}{r^2} = \oint_a d\mathbf{l}_a \oint_b \frac{d\mathbf{l}_b \cdot \hat{\mathbf{r}}}{r^2}.$$

- This integral is zero because b is closed and

$$\mathbf{F}_{ab} = -\frac{\mu_0}{4\pi} I_a I_b \oint_a \oint_b \hat{\mathbf{r}} \frac{d\mathbf{l}_a \cdot d\mathbf{l}_b}{r^2}. \quad (22-23)$$

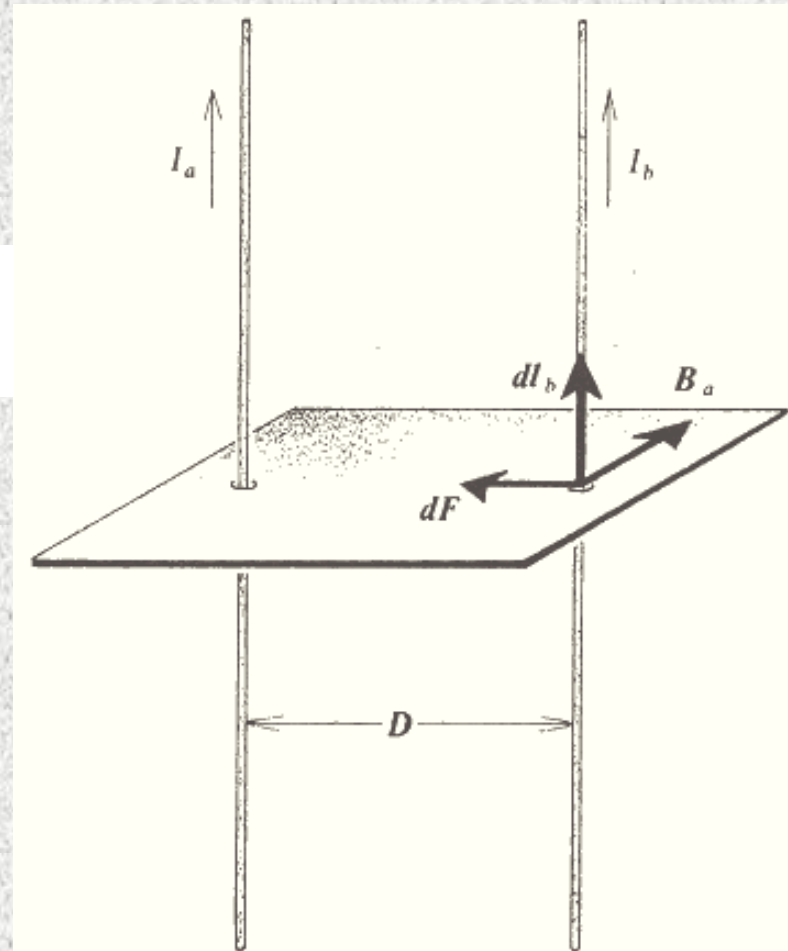


Force between two parallel Currents

- Without integration, we can calculate the force per unit length between two long parallel wires bearing current
- At the position of I_b , B_a is $\mu_0 I_a / 2\pi D$ in the direction shown in the figure
- The force on a unit length of wire b is thus

$$F' = B_a I_b = \frac{\mu_0 I_a I_b}{2\pi D}. \quad (22-24)$$

- The force is attractive if the current flow in the same direction and repulsive otherwise.
- This force is normally negligible



Magnetic Force on a volume Distribution of Current

- Consider a small element of volume of length $d\mathbf{l}$ parallel to \mathbf{J} and of cross section of dA , as in figure.

- The magnetic force on the element is

$$d\mathbf{F} = (\mathbf{J} d\mathcal{A}) d\mathbf{l} \times \mathbf{B} = \mathbf{J} \times \mathbf{B} dv$$

- The force per unit volume is

$$\mathbf{F}' = \mathbf{J} \times \mathbf{B}. \quad (22-30)$$

- The total magnetic force on a given distribution of conducting currents occupying a volume v is

$$\mathbf{F} = \int_v \mathbf{J} \times \mathbf{B} dv. \quad (22-31)$$

